

Lecture 21

• Stokes' theorem

Green's theorem states that

$$\iint_D \left(\frac{\partial Q}{\partial x} - \frac{\partial P}{\partial y} \right) dA = \int_C P dx + Q dy.$$

Stokes' theorem is the curved version of Green's theorem when the plane region D is replaced by a surface S in \mathbb{R}^3 and the curve C is replaced by a curve C in space.

To guess what Stokes' theorem would look like, consider

$$\vec{F} = P \hat{i} + Q \hat{j} + R \hat{k}.$$

When C is a circle lying on the x - y plane, one would get the formula above.

When C is a circle on the y - z plane, we get

$$\iint_S \left(\frac{\partial R}{\partial y} - \frac{\partial Q}{\partial z} \right) dA = \int_C Q dy + R dz$$

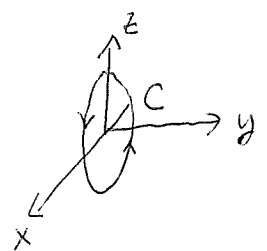
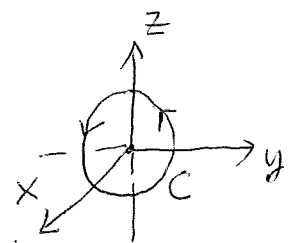
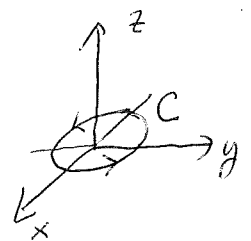
When C is a circle on the z - x plane, we get

$$\iint_S \left(\frac{\partial P}{\partial z} - \frac{\partial R}{\partial x} \right) dA = \int_C P dx + R dz.$$

Therefore, for a general S lying in space,

the RHS of Stokes' should contain

$$\frac{\partial Q}{\partial x} - \frac{\partial P}{\partial y}, \frac{\partial R}{\partial y} - \frac{\partial Q}{\partial z}, \frac{\partial P}{\partial z} - \frac{\partial R}{\partial x}.$$



Let $\vec{F} = (P, Q, R)$ or $P\hat{i} + Q\hat{j} + R\hat{k}$ a C^1 -v.f. in \mathbb{R}^3 . Define its curl to be a v.f.

$$\text{curl } \vec{F} = \left(\frac{\partial R}{\partial y} - \frac{\partial Q}{\partial z}\right)\hat{i} + \left(-\frac{\partial R}{\partial x} + \frac{\partial P}{\partial z}\right)\hat{j} + \left(\frac{\partial Q}{\partial x} - \frac{\partial P}{\partial y}\right)\hat{k}.$$

Let $\nabla = (\partial_x, \partial_y, \partial_z)$ be the "gradient operator". The curl vector could be computed via

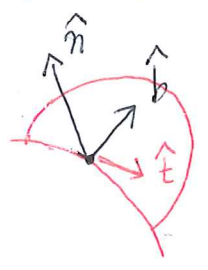
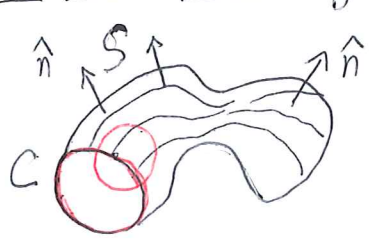
$$\det \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ \partial_x & \partial_y & \partial_z \\ P & Q & R \end{vmatrix}.$$

The notation $\nabla \times \vec{F}$ is often used instead of $\text{curl } \vec{F}$.

Theorem 1 (Stokes' thm) Let S be an oriented surface $\subset \mathbb{R}^3$ whose boundary is a C^1 simple, closed curve C . For a C^1 -v.f. \vec{F} on S ,

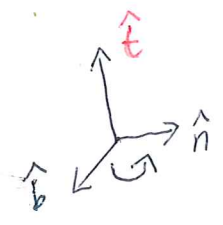
$$\iint_S \nabla \times \vec{F} \cdot \hat{n} \, d\sigma = \oint_C P dx + Q dy + R dz.$$

Here $\vec{F} = P\hat{i} + Q\hat{j} + R\hat{k}$ and \hat{n} is the chosen unit normal on S . The orientation of C is determined by \hat{n} according to the following right hand rule.



\hat{b} = binormal, the ^{tangent} vector perpendicular to the boundary and pointing inward S .

$$\hat{b} \rightarrow \hat{n} \rightarrow \hat{t}$$



If we take S lying on the x - y plane and $\hat{n} = \hat{k}$, the boundary C is on the xy -plane running in anticlockwise direction. So

$$\begin{aligned} \nabla \times \vec{F} \cdot \hat{n} &= \nabla \times \vec{F} \cdot \hat{k} \\ &= \frac{\partial Q}{\partial x} - \frac{\partial P}{\partial y}. \end{aligned}$$

Any parametrization of C is of the form $(x(t), y(t), 0)$ so $\vec{r}'(t) = (x'(t), y'(t), 0)$. Thus

$$\begin{aligned} &\oint_C P dx + Q dy + R dz \\ &= \int_a^b (P(\vec{r}(t)) x'(t) + Q(\vec{r}(t)) y'(t) + R(\vec{r}(t)) 0) dt \\ &= \int_a^b P(\vec{r}(t)) x'(t) + Q(\vec{r}(t)) y'(t) dt \\ &= \oint_C P dx + Q dy. \end{aligned}$$

So,
$$\iint_S \nabla \times \vec{F} \cdot \hat{n} d\sigma = \iint_S \left(\frac{\partial Q}{\partial x} - \frac{\partial P}{\partial y} \right) dA,$$

$$\oint_C P dx + Q dy + R dz = \oint_C P dx + Q dy,$$

Stokes' thm reduces to Green's Theorem.

Proof of Stokes' thm. We give a proof in the special case where

$$S = \{ (x, y, f(x, y)) : (x, y) \in D \}$$

and D is bounded by a simple, closed curve C_1 .

The standard parametrization of S

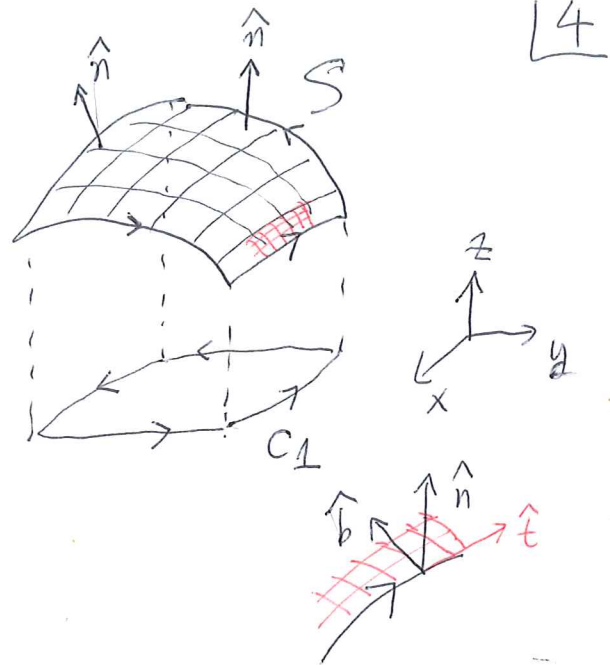
$$(x, y) \mapsto (x, y, f(x, y))$$

$$\vec{r}_x = (1, 0, f_x)$$

$$\vec{r}_y = (0, 1, f_y)$$

$$\vec{r}_x \times \vec{r}_y = (-f_x, -f_y, 1)$$

$$\hat{n} = \frac{(-f_x, -f_y, 1)}{\sqrt{1 + |\nabla f|^2}}$$



$$\therefore \iint_S \nabla \times \vec{F} \cdot \hat{n} \, d\sigma = \iint_S \nabla \times \vec{F} \cdot \frac{(-f_x, -f_y, 1)}{\sqrt{1 + |\nabla f|^2}} \, d\sigma$$

$$= \iint_D \nabla \times \vec{F} \cdot (-f_x, -f_y, 1) \, dA(x, y)$$

$$= \iint_D [(R_y - Q_z)(-f_x) + (-R_x + P_z)(-f_y) + (Q_x - P_y)] \, dA(x, y) \quad \text{--- (1)}$$

On the other hand, let $\vec{\gamma}(t) = (x(t), y(t), f(x(t), y(t)))$ be a parametrization of C , the boundary of S , (so $t \mapsto (x(t), y(t))$ parametrizes C_1).

$$\begin{aligned} \oint_C P \, dx + Q \, dy + R \, dz &= \int_a^b [P(\vec{\gamma}(t)) x'(t) + Q(\vec{\gamma}(t)) y'(t) + R(\vec{\gamma}(t)) z'(t)] \, dt \\ &= \int_a^b P(\vec{\gamma}(t)) x'(t) + Q(\vec{\gamma}(t)) y'(t) + R(\vec{\gamma}(t)) (f_x x' + f_y y') \, dt \\ &= \int_a^b [(P + R f_x) x' + (Q + R f_y) y'] \, dt \end{aligned}$$

note that $P + R f_x = P(\vec{\gamma}(t)) + R(\vec{\gamma}(t)) f_x(x(t), y(t))$, etc.

Let $\tilde{P}(x, y) = P(x, y, f(x, y)) + R(x, y, f(x, y))f_x(x, y)$, &
 $\tilde{Q}(x, y) = Q(x, y, f(x, y)) + R(x, y, f(x, y))f_y(x, y)$.

Then $\int_a^b [(P + Rf_x)x' + (Q + Rf_y)y'] dt$

$= \oint_{C_1} \tilde{P} dx + \tilde{Q} dy$

(apply Green's to $\tilde{P} \hat{i} + \tilde{Q} \hat{j}$ on D)

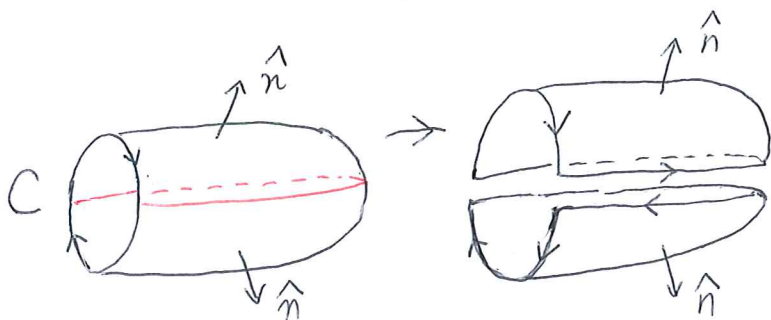
$= \iint_D \left(\frac{\partial \tilde{Q}}{\partial x} - \frac{\partial \tilde{P}}{\partial y} \right) dA$

$= \iint_D [Q_x + Q_z f_x + \cancel{(R_x + R_z f_x)} f_y + R f_{yx} - P_y - P_z f_y - \cancel{(R_y + R_z f_y)} f_x - \cancel{R f_{xy}}] dA$

$= \iint_D [(Q_z - R_y) f_x + (R_x - P_z) f_y + Q_x - P_y] dA$

$= \textcircled{1}$, done.

As in the case of Green's thm, the general case can be proved by cutting up S .



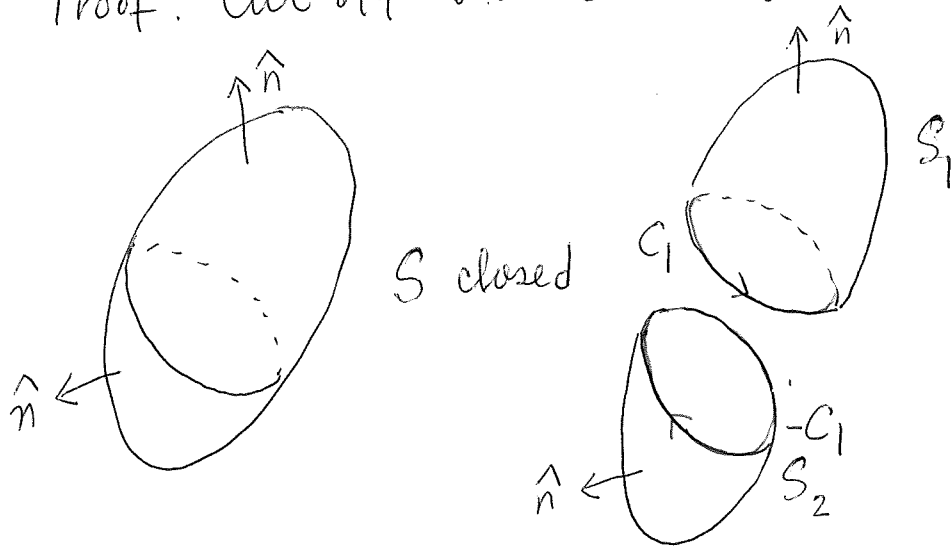
$S \rightarrow S_1, S_2$
 both S_1, S_2 are graphs over xy -plane.

Corollary 1

when S is a closed surface, that's, C is empty,
Stokes' Theorem becomes

$$\oiint_S \nabla \times \vec{F} \cdot \hat{n} \, d\sigma = 0$$

Proof. Cut off into two, to induce C and $-C$.



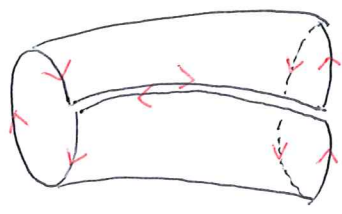
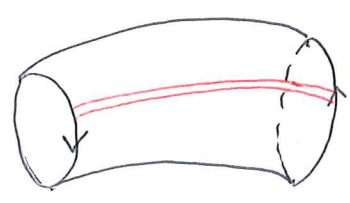
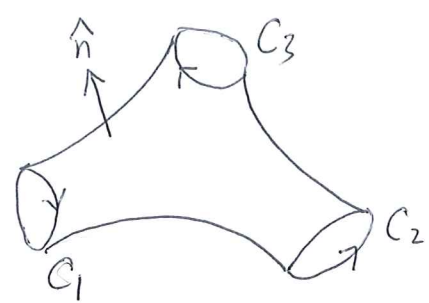
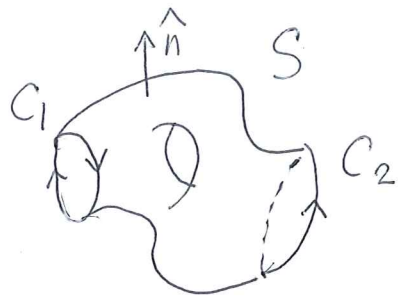
$$\begin{aligned} \oiint_S &= \iint_{S_1} + \iint_{S_2} \\ &= \oint_{C_1} + \oint_{-C_1} = 0. \end{aligned}$$

Corollary 2

when S bounds several simple, closed curves C_1, \dots, C_n ,

$$\oiint_S \nabla \times \vec{F} \cdot \hat{n} \, d\sigma = \sum_{j=1}^n \oint_{C_j} \vec{F} \cdot d\vec{r}$$

Pf: Cutting



a single long boundary C

e.g. 1 Let S be the upper hemisphere $x^2 + y^2 + z^2 = 1, z \geq 0$, with normal pointing out. Evaluate

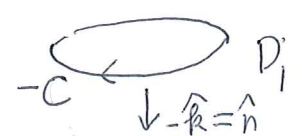
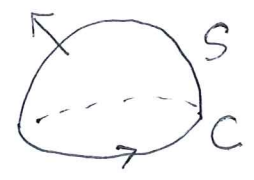
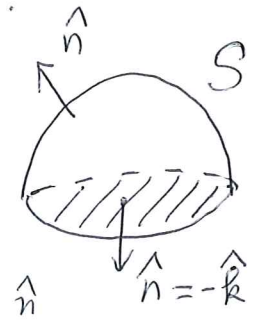
$$\iint_S \nabla \times \vec{F} \cdot \hat{n} \, d\sigma, \quad \vec{F} = y\hat{i} + x^2z\hat{j} + e^{xz}\hat{k}$$

We have

$$\nabla \times \vec{F} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ \partial_x & \partial_y & \partial_z \\ y & x^2z & e^{xz} \end{vmatrix}$$

$$= -x\hat{i} - ze^{xz}\hat{j} + (2xz - 1)\hat{k}$$

Let $S_1 = D_1$ the x - y plane unit disk sitting in \mathbb{R}^3 as a surface with normal pointing down (so that $S \cup S_1$ becomes a closed surface with outer normal).



By $\iint_{S \cup S_1} \nabla \times \vec{F} \cdot \hat{n} \, d\sigma = 0,$

$$\iint_S \nabla \times \vec{F} \cdot \hat{n} \, d\sigma = - \iint_{S_1} \nabla \times \vec{F} \cdot \hat{n} \, d\sigma$$

$$= - \iint_{S_1} \left(x^2 \hat{i} - z e^{xz} \hat{j} + (2xz-1) \hat{k} \right) \cdot \left(-\hat{k} \right) d\sigma$$

$$= \iint_{S_1} (2xz-1) d\sigma$$

$$d\sigma = \sqrt{1 + |\nabla f|^2} \, dA$$

$$z = f(x,y) \equiv 0$$

$$d\sigma = dA$$

Using the parametrization $(x, y) \mapsto (x, y, 0)$,

$$= \iint_{D_1} -dA$$

$$= -\pi \cdot \#$$

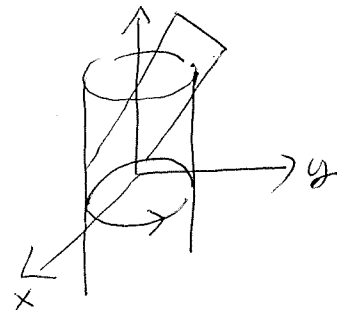
e.g. 2 Evaluate

$$\oint_C (y+z) dx + (z+x) dy + x dz \quad \text{where}$$

C is the intersection of $x^2 + y^2 = 1$ (cylinder)

and $x+y+z=10$ (plane).

the orientation is anticlockwise direction.



$$\nabla \times \vec{F} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ \partial_x & \partial_y & \partial_z \\ y+z & z+x & x \end{vmatrix}$$

$$= -\hat{i} + 0\hat{j} + 0\hat{k} = -\hat{i}$$

$(x, y) \mapsto (x, y, \frac{1}{2}(10-x-y))$ parametrise the plane,

$$\hat{n} = \frac{(-f_x - f_y, 1)}{\sqrt{1 + |\nabla f|^2}} = \frac{(\frac{1}{2}, \frac{1}{2}, 1)}{\sqrt{\frac{3}{2}}} \text{ pointing up.}$$

Using Green's theorem, we have shown that the compatibility condition $P_y = Q_x$ is also sufficient when G is simply-connected.

Theorem ($n=2$) As above, in addition G is simply connected.

then $\frac{\partial P}{\partial y} = \frac{\partial Q}{\partial x} \Rightarrow \vec{F} = P\hat{i} + Q\hat{j}$ has a potential.

Now, we have

Theorem ($n=3$) As above, in addition G is simply connected.

then

$$\frac{\partial Q}{\partial z} = \frac{\partial R}{\partial y}, \quad \frac{\partial P}{\partial z} = \frac{\partial R}{\partial x}, \quad \frac{\partial P}{\partial y} = \frac{\partial Q}{\partial x} \quad \text{or} \quad \nabla \times \vec{F} = \vec{0}.$$

$$\Rightarrow \vec{F} = P\hat{i} + Q\hat{j} + R\hat{k} \quad \text{has a potential.}$$

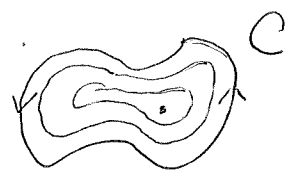
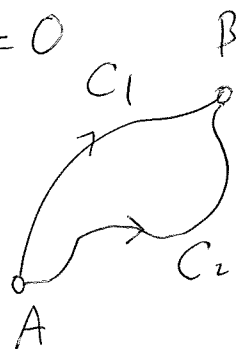
PF: We first show that

$$\oint_C \vec{F} \cdot d\vec{r} = \oint_C P dx + Q dy + R dz = 0$$

for every simple, closed curve in G .

Using this property we can define

the potential of \vec{F} , see



By assumption, G is simply connected. We can deform C into a point. The deformation forms a surface whose

boundary is C . Let the orientation of S be induced by the orientation of C . By Stokes' theorem,

$$\oint_C \vec{F} \cdot d\vec{r} = \iint_S \nabla \times \vec{F} \cdot \hat{n} \, d\sigma.$$

However, one observes that the compatibility conditions are nothing but $\nabla \times \vec{F} = \vec{0}$. So

$$\oint_C \vec{F} \cdot d\vec{r} = 0, \quad \text{done. } \#$$

So, finally we have shown that in a simply connected region, a vector field is conservative if and only if

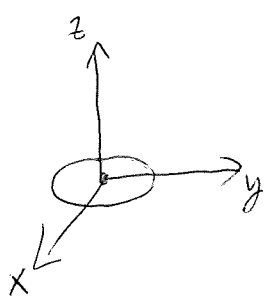
$$\nabla \times \vec{F} = \vec{0},$$

that is, it is curl free. We can express this result

as $\text{curl } \vec{F} = 0 \iff \exists \Phi \text{ s.t. } \vec{F} = \nabla \Phi.$

Note that $\nabla \times \nabla \Phi = 0$ (a conservative v.f. has no curl) no matter the region is simply connected or not.

Next, we examine the meaning of $\nabla \times \vec{F}$.



Let (x_0, y_0, z_0) be a point and use it as the origin.

Consider the circle $C_\epsilon = (\epsilon \cos \theta, \epsilon \sin \theta, 0)$
 $\theta \in [0, 2\pi]$.

Stokes theorem :

$$\oint_{C_\epsilon} P dx + Q dy + R dz = \iint_{D_\epsilon} \nabla \times \vec{F} \cdot \hat{k} d\sigma$$

$$\therefore \lim_{\epsilon \rightarrow 0} \frac{1}{|D_\epsilon|} \oint_{C_\epsilon} P dx + Q dy + R dz$$

circulation around
 C_ϵ

$$= \lim_{\epsilon \rightarrow 0} \frac{1}{|D_\epsilon|} \iint_{D_\epsilon} \left(\frac{\partial Q}{\partial x} - \frac{\partial P}{\partial y} \right) dA$$

$$= \left(\frac{\partial Q}{\partial x} - \frac{\partial P}{\partial y} \right) (x_0, y_0, z_0)$$

So $\left(\frac{\partial Q}{\partial x} - \frac{\partial P}{\partial y} \right) (x_0, y_0, z_0)$ measures the strength of the v.f. \vec{F} rotates around the z-axis.

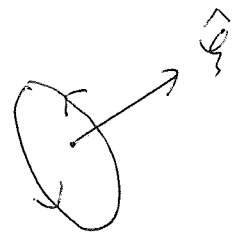
Similarly,

$\left(\frac{\partial R}{\partial y} - \frac{\partial Q}{\partial z} \right) (x_0, y_0, z_0)$ measures the strength of \vec{F} at (x_0, y_0, z_0) when rotates around the x-axis.

$\left(-\frac{\partial R}{\partial x} + \frac{\partial P}{\partial z} \right) (x_0, y_0, z_0)$ measure -----

----- y-axis.

How about around the z-axis?

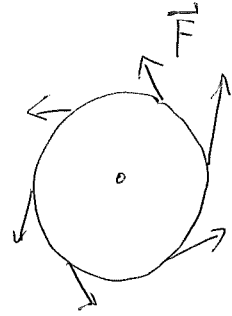


It is

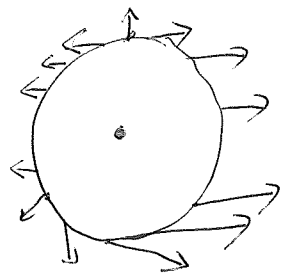
$$\lim_{\epsilon \rightarrow 0} \frac{1}{|D_\epsilon|} \oint_{C_\epsilon} P dx + Q dy + R dz$$

$$= \lim_{\epsilon \rightarrow 0} \frac{1}{|D_\epsilon|} \iint_{D_\epsilon} \nabla \times \vec{F} \cdot \hat{\xi} d\sigma$$

$$= \nabla \times \vec{F} \cdot \hat{\xi} \quad \#$$



when \vec{F} rotates around the pt, the curl at the pt > 0 or < 0 .



\vec{F} not nec. around the pt, curl \vec{F} could still be non-zero.